

# Optomechanical cooling efficiency: the cost of turning a valve

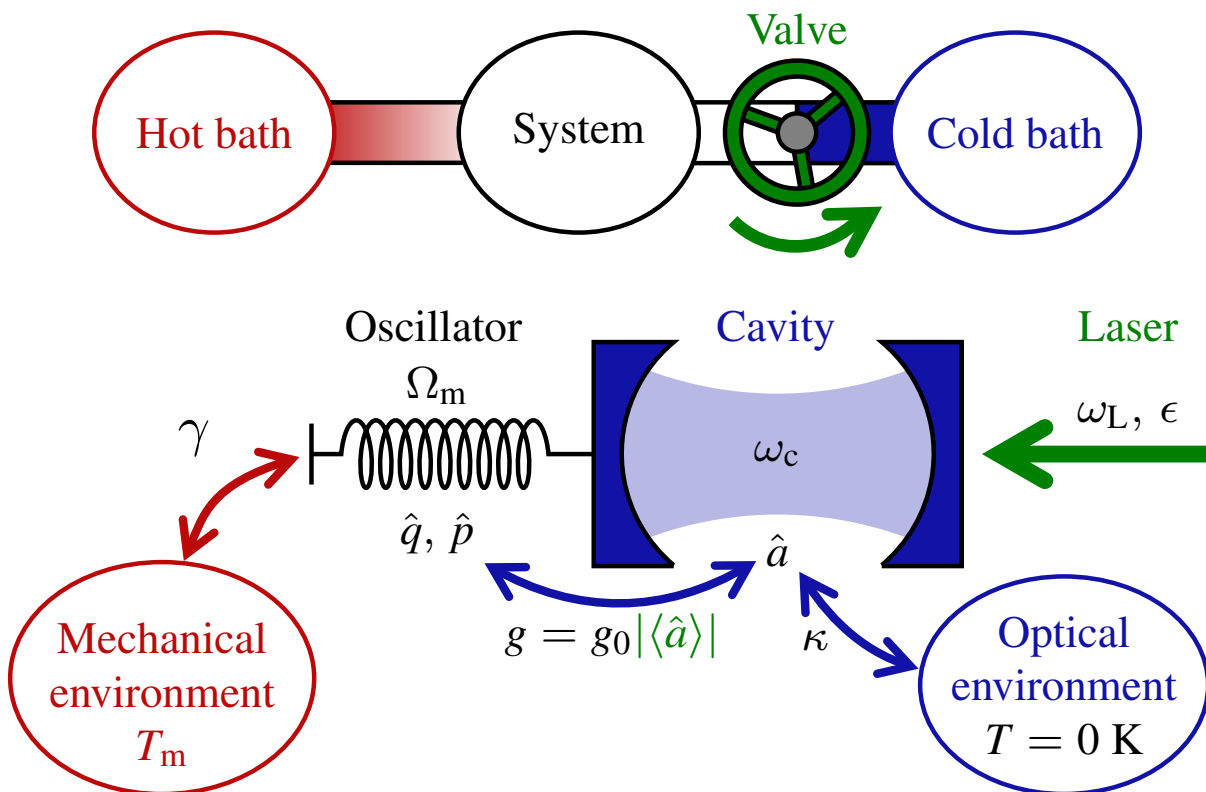


CHALMERS  
UNIVERSITY OF TECHNOLOGY

Juliette Monsel<sup>a\*</sup>, Nastaran Dashti<sup>a</sup>, Sushanth Kini Manjeshwar<sup>a</sup>, Janine Splettstoesser<sup>a</sup> and Witlef Wieczorek<sup>a</sup>

<sup>a</sup> Department of Microtechnology and Nanoscience (MC2), Chalmers University of Technology, S-412 96 Göteborg, Sweden

\* monsel@chalmers.se



**Platform:** Cavity optomechanics

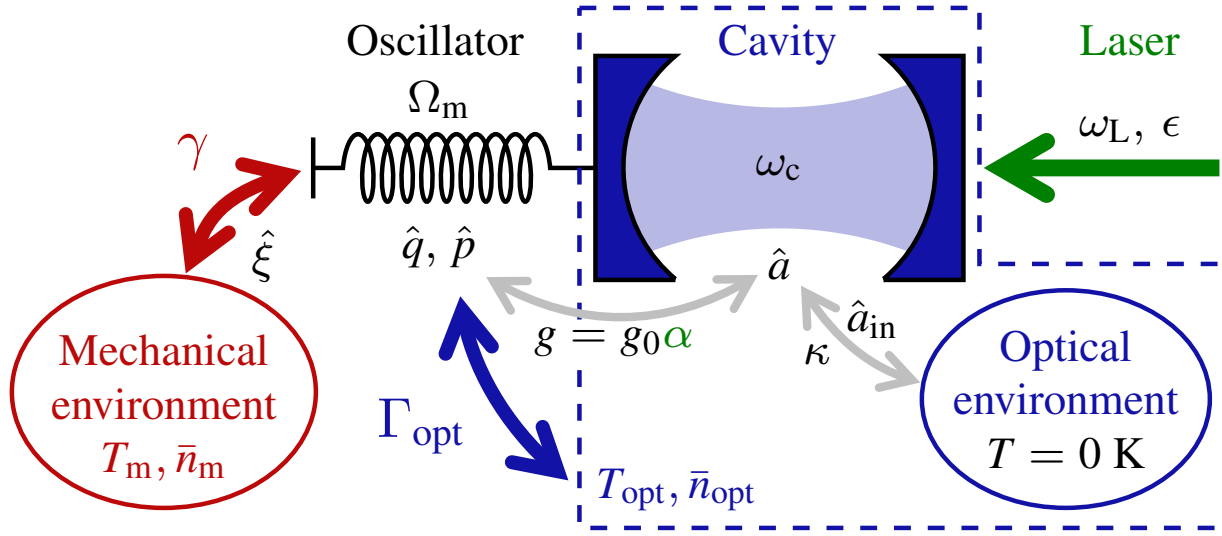
**Protocol:** Cooling down the mechanical oscillator in the resolved-sideband regime ( $\kappa \ll \Omega_m$ )

- ▶ How can we understand resolved-sideband cooling from a thermodynamic perspective?

Improved cooling schemes including using squeezed light or a Fano mirror have been proposed:

- ▶ How do they compare?
- ▶ Are they more efficient than the original scheme?

# 1. Resolved-sideband cooling



$$H = \hbar \Delta \hat{a}^\dagger \hat{a} + \frac{\hbar \Omega_m}{2} (\hat{q}^2 + \hat{p}^2) - \hbar g_0 \sqrt{2} \hat{q} \hat{a}^\dagger \hat{a} + \hbar (\epsilon \hat{a}^\dagger + \epsilon^* \hat{a})$$

Linearization<sup>1</sup>:

$$\hat{a} = \alpha + \delta \hat{a}$$

$$\hat{q} = \bar{q} + \delta \hat{q}$$

$$\hat{p} = \bar{p} + \delta \hat{p}$$

$$\alpha = -i\epsilon / (\kappa + i\tilde{\Delta})$$

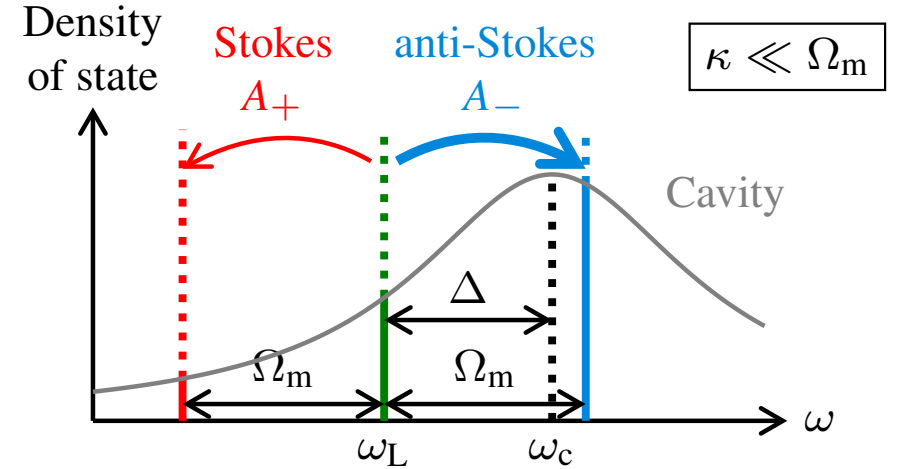
$$\bar{q} = \sqrt{2} g_0 |\alpha|^2 / \Omega_m$$

$$\bar{p} = 0$$

$$\tilde{\Delta} = \omega_c - \omega_L - \sqrt{2} g_0 \bar{q}$$

$$H = \hbar \tilde{\Delta} \delta \hat{a}^\dagger \delta \hat{a} + \frac{\hbar \Omega_m}{2} (\delta \hat{q}^2 + \delta \hat{p}^2) - \hbar g \sqrt{2} \delta \hat{q} (\delta \hat{a} + \delta \hat{a}^\dagger)$$

1. C. Genes et al. *Phys. Rev. A* **77**, 033804 (2008)



Scattering processes:

► **Stokes:** photon  $\omega_L \rightarrow$  phonon  $\Omega_m$  + photon  $\omega_L - \Omega_m$

► **anti-Stokes:** photon  $\omega_L$  + phonon  $\Omega_m \rightarrow$  photon  $\omega_L + \Omega_m$

Dynamics:

$$\delta \dot{\hat{q}} = \Omega_m \delta \hat{p}$$

Langevin equations<sup>1</sup>

$$\delta \dot{\hat{p}} = -\Omega_m \delta \hat{q} - \gamma \delta \hat{p} + g \sqrt{2} (\delta \hat{a}^\dagger + \delta \hat{a}) + \sqrt{\gamma} \hat{\xi}$$

$$\delta \dot{\hat{a}} = -(\kappa + i\tilde{\Delta}) \delta \hat{a} + ig \sqrt{2} \delta \hat{q} + \sqrt{2\kappa} \hat{a}_{in}$$

Noises

$$\langle \hat{\xi}(t) \hat{\xi}(t') \rangle = (2\bar{n} + 1) \delta(t - t')$$

$$\langle \hat{a}_{in}(t) \hat{a}_{in}^\dagger(t') \rangle = \delta(t - t')$$

## 2. Thermodynamic analysis

**Weak coupling regime** ( $g \ll \kappa, \Omega_m$ ):

Steady state phonon number:

$$n_{\text{eff}} = \langle \delta \hat{b}^\dagger \delta \hat{b} \rangle_{\text{ss}} = \frac{\gamma \bar{n}_m + \Gamma_{\text{opt}} \bar{n}_{\text{opt}}}{\gamma + \Gamma_{\text{opt}}}$$

with  $\Gamma_{\text{opt}} = A_- - A_+$ ,  $\bar{n}_{\text{opt}} = A_+ / \Gamma_{\text{opt}}$

⇒ Optical part behaves like a cold phonon bath at effective temperature  $T_{\text{opt}}$  such that  $\bar{n}_{\text{opt}} = (e^{\hbar\Omega_m/k_B T_{\text{opt}}} - 1)^{-1}$

1. I. Wilson-Rae et al. *New J. Phys.* **10**, 095007 (2008)

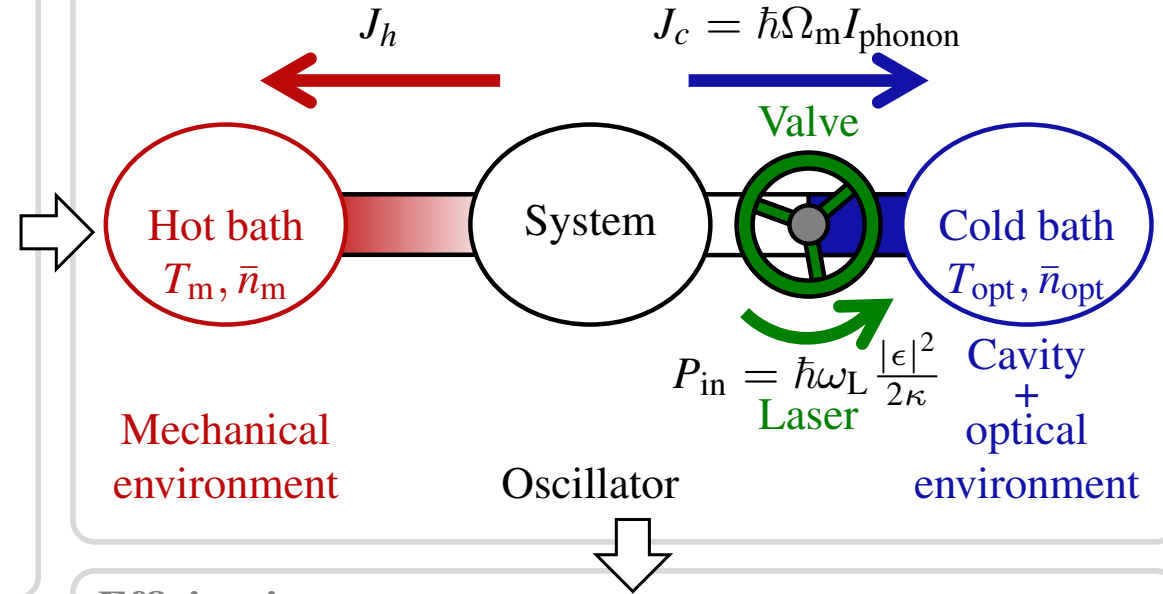
**Beyond the weak coupling regime:**

$$\frac{d\langle \delta \hat{a}^\dagger \delta \hat{a} \rangle}{dt} = ig \left( \langle \delta \hat{a}^\dagger (\delta \hat{b}^\dagger + \delta \hat{b}) \rangle - \langle \delta \hat{a} (\delta \hat{b}^\dagger + \delta \hat{b}) \rangle \right) \leftarrow I_{\text{photon}} - 2\kappa \langle \delta \hat{a}^\dagger \delta \hat{a} \rangle$$

$$\begin{aligned} \frac{d\langle \delta \hat{b}^\dagger \delta \hat{b} \rangle}{dt} = & ig \left( \langle \delta \hat{b}^\dagger (\delta \hat{a} + \delta \hat{a}^\dagger) \rangle - \langle \delta \hat{b} (\delta \hat{a}^\dagger + \delta \hat{a}) \rangle \right) \leftarrow I_{\text{photon}} \\ & + \gamma (\bar{n} - \langle \delta \hat{b}^\dagger \delta \hat{b} \rangle) + \frac{\gamma}{2} \left( 1 + \langle \delta \hat{b}^2 \rangle + \langle (\delta \hat{b}^\dagger)^2 \rangle \right) \end{aligned}$$

... and the equations for all the other 2nd order moments ⇒ steady-state

**Thermodynamic roles:**



**Efficiencies:**

- ▶ Resource = flux of photons interacting with the mechanics

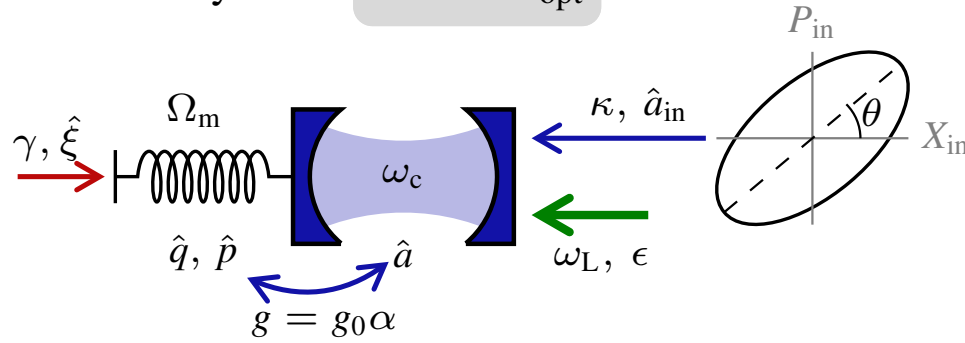
$$\eta_{\text{conv}} = \frac{\hbar\Omega_m I_{\text{photon}}}{\hbar\tilde{\Delta} I_{\text{photon}}}$$

- ▶ Resource = laser power

$$\eta_{\text{laser}} = \frac{\hbar\Omega_m I_{\text{photon}}}{P_{\text{in}}}$$

### 3. Improved cooling schemes

**Squeezed light:** Squeeze the input light to decrease the amplitude noise in the cavity<sup>1,2</sup>  $\Rightarrow$  Lowers  $T_{\text{opt}}$



Cavity's input noise:  $\langle \hat{a}_{\text{in}}^\dagger(t) \hat{a}_{\text{in}}(t') \rangle = N \delta(t - t')$

$\langle \hat{a}_{\text{in}}(t) \hat{a}_{\text{in}}(t') \rangle = M \delta(t - t')$

with

$$N = \frac{4r_s^2}{(1 - r_s^2)^2}, \quad M = \frac{2r_s(1 + r_s^2)}{(1 - r_s^2)^2} e^{-i2\theta},$$

and squeezing ratio  $0 < r_s < 1$ .

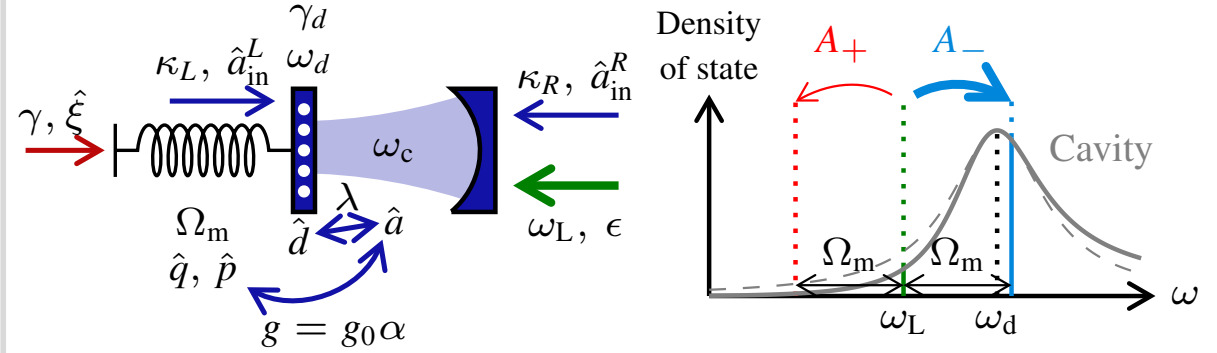
$\Rightarrow$  Optimal  $\theta$  and  $r_s$  minimizes amplitude noise

$\Rightarrow$  **Cost:** Squeezing generation

1. M. Asjad et al. *Phys. Rev. A* **94**, 051801 (2016)

2. J. B. Clark et al. *Nature* **541**, 191 (2017)

**Fano mirror:** Asymmetrize the cavity's lineshape to suppress the Stokes process<sup>3</sup>  $\Rightarrow$  Lowers  $T_{\text{opt}}$  and increases  $\Gamma_{\text{opt}}$



$\delta \dot{\hat{q}} = \Omega_m \delta \hat{p}$ ,

$\delta \dot{\hat{p}} = -\Omega_m \delta \hat{q} - \gamma \delta \hat{p} + g\sqrt{2}(\delta \hat{a}^\dagger + \delta \hat{a}) + \sqrt{\gamma} \hat{\xi}$

$\delta \dot{\hat{a}} = -(\kappa + i\tilde{\Delta})\delta \hat{a} + ig\sqrt{2}\delta \hat{q} - \mathcal{G}\delta \hat{d} + \sqrt{2\kappa_L}\hat{a}_{\text{in}}^L + \sqrt{2\kappa_R}\hat{a}_{\text{in}}^R$

$\delta \dot{\hat{d}} = -(\gamma_d + i\tilde{\Delta}_d)\delta \hat{d} - \mathcal{G}\delta \hat{a} + \sqrt{2\gamma_d}\hat{a}_{\text{in}}^L$

with  $\mathcal{G} = i\lambda + \sqrt{\kappa_L\gamma_d}$ ,  $\kappa = \kappa_L + \kappa_R$

$\Rightarrow$  Optimal  $\gamma_d$  simultaneously maximize cooling and minimizing heating for  $\tilde{\Delta}_d \simeq \Omega_m$

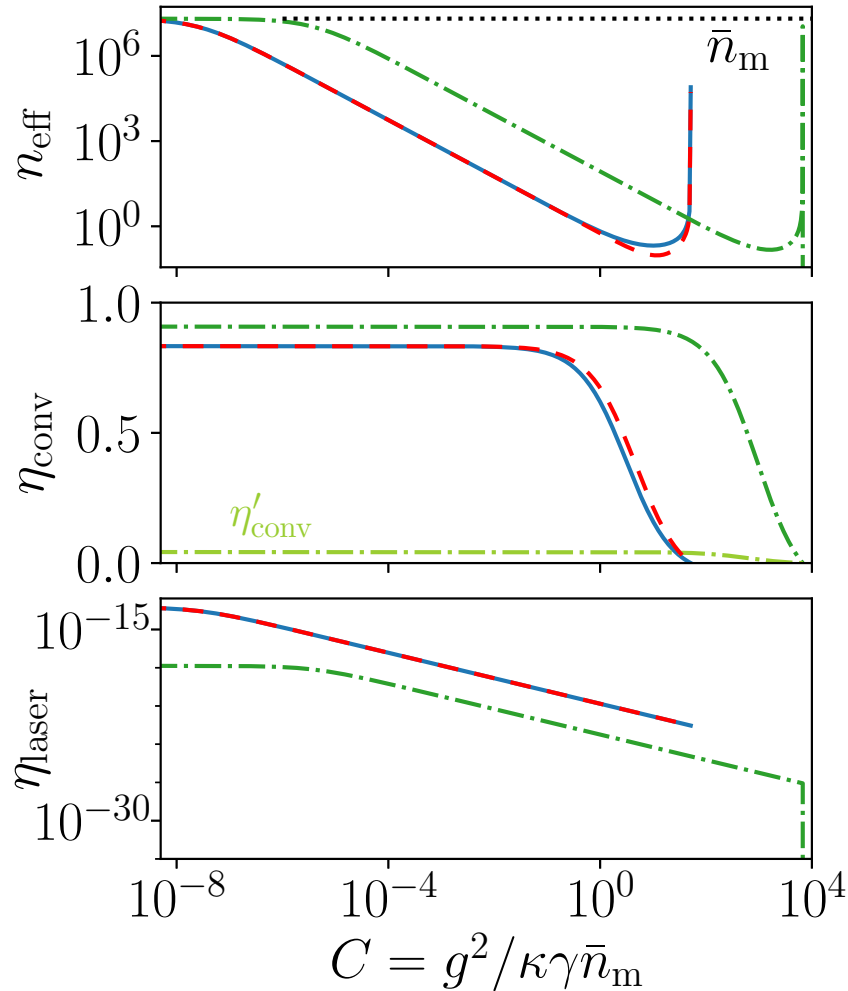
$\Rightarrow$  **Cost:** Build field in Fano mirror's mode

$$\eta'_{\text{conv}} = \frac{\Omega_m g \text{Im}(\langle \delta \hat{b}^\dagger (\delta \hat{a}^\dagger + \delta \hat{a}) \rangle)}{\tilde{\Delta} (g \text{Im}(\langle \delta \hat{a} (\delta \hat{b}^\dagger + \delta \hat{b}) \rangle) - \text{Re}(\mathcal{G} \langle \delta \hat{d} \delta \hat{a}^\dagger \rangle))}$$

3. O. Černotík et al. *Phys. Rev. Lett.* **122**, 243601 (2019)

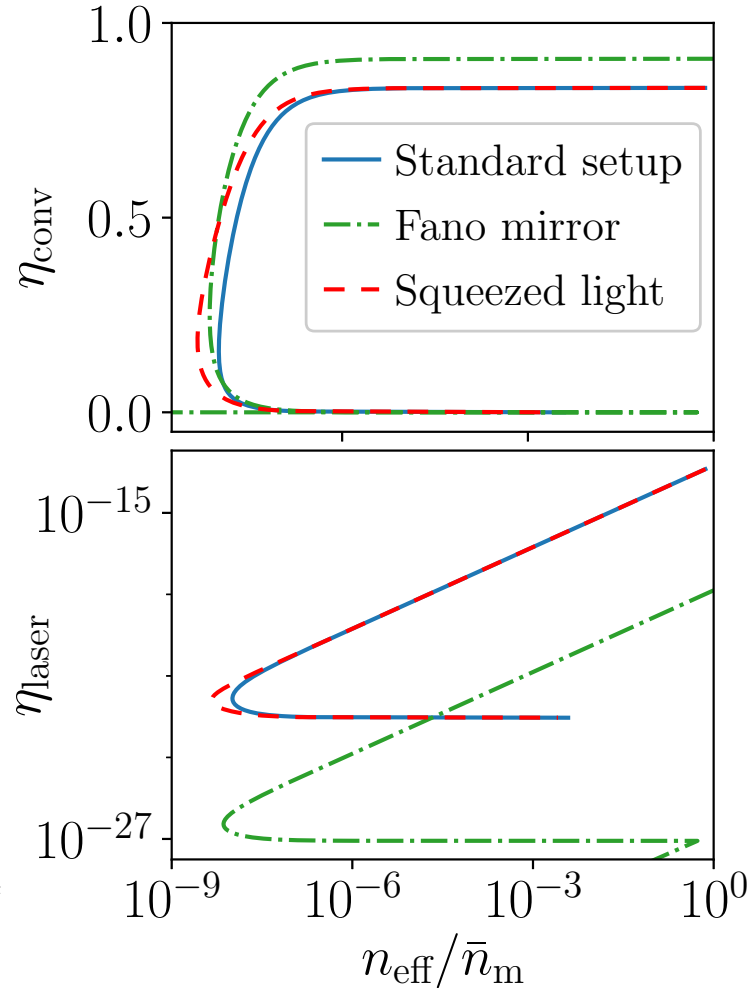
## 4. Phonon number and efficiencies

As functions of cooperativity  $C$



$$\eta_{\text{conv}} = \frac{\hbar \Omega_m I_{\text{phonon}}}{\hbar \tilde{\Delta} I_{\text{photon}}} \quad \text{and} \quad \eta_{\text{laser}} = \frac{\hbar \Omega_m I_{\text{phonon}}}{\hbar \omega_L |\epsilon|^2 / 2\kappa}$$

Efficiency versus phonon number



**Parameters:** (Levitated optomechanics)

$$\begin{aligned}
 \Omega_m / 2\pi &= 305 \text{ kHz} & \lambda_L &= 1064 \text{ nm} \\
 T_m &= 300 \text{ K} & \tilde{\Delta} &= \Omega_m \\
 \gamma / 2\pi &= 1.6 \cdot 10^{-4} \text{ Hz} & \kappa / 2\pi &= 193 \text{ kHz} \\
 \bar{n}_m &= 1.3 \cdot 10^8 & \kappa / \Omega_m &= 0.63 \\
 Q_m &= 1.9 \cdot 10^9 & g_0 / 2\pi &= 0.3 \text{ kHz} \\
 C_0 &= 2.2 \cdot 10^{-11}
 \end{aligned}$$

U. Delić et al. *Science* **367**, 892 (2020)

Squeezing:  $r_s = 0.16$ ,  $\theta = 0.89$  rad (optimal)

Fano mirror:  $\tilde{\Delta}_d = \Omega_m$ ,  $\gamma_d / 2\pi = 3.86$  MHz (optimal), other parameters chosen to have an effective cavity linewidth =  $\kappa$

**Squeezed light:** Smallest  $n_{\text{eff}}$

- ⇒ Efficiencies barely impacted
- ⇒ Cost of squeezing generation not taken into account

**Fano mirror:** Smaller  $n_{\text{eff}}$  (but  $> n_{\text{eff}}^{\text{squeezing}}$ )

- ⇒ Significant decrease in efficiency
- ⇒ Cost of building a coherent field in the Fano mirror's mode

## 5. Conclusion

- ▶ Cavity + Optical environment = Effective cold phonon bath
- ▶ Optomechanical coupling  $g = g_0\alpha$  is like a valve  
⇒ Cost to couple the cold bath to the system: Build coherent field  $\alpha$  in the cavity
- ▶ Combining dynamics, beyond the weak coupling regime, with thermodynamic analysis  
⇒ Efficiencies  $\eta_{\text{conv}}$  (conversion photon  $\rightarrow$  phonon) and  $\eta_{\text{laser}}$  (resource = laser power)